Gravitino Dark Matter and Related Issues

The CMSSM after WMAP Detectability B_s → μ⁺ μ⁻
The VCMSSM (mSUGRA) Gravitino Dark Matter BBN limits - Li problem

with: Ellis, Santoso, Spanos

Boundary conditions

• Input parameters: μ , m_1 , m_2 , B. predict M_Z , tan β , m_A

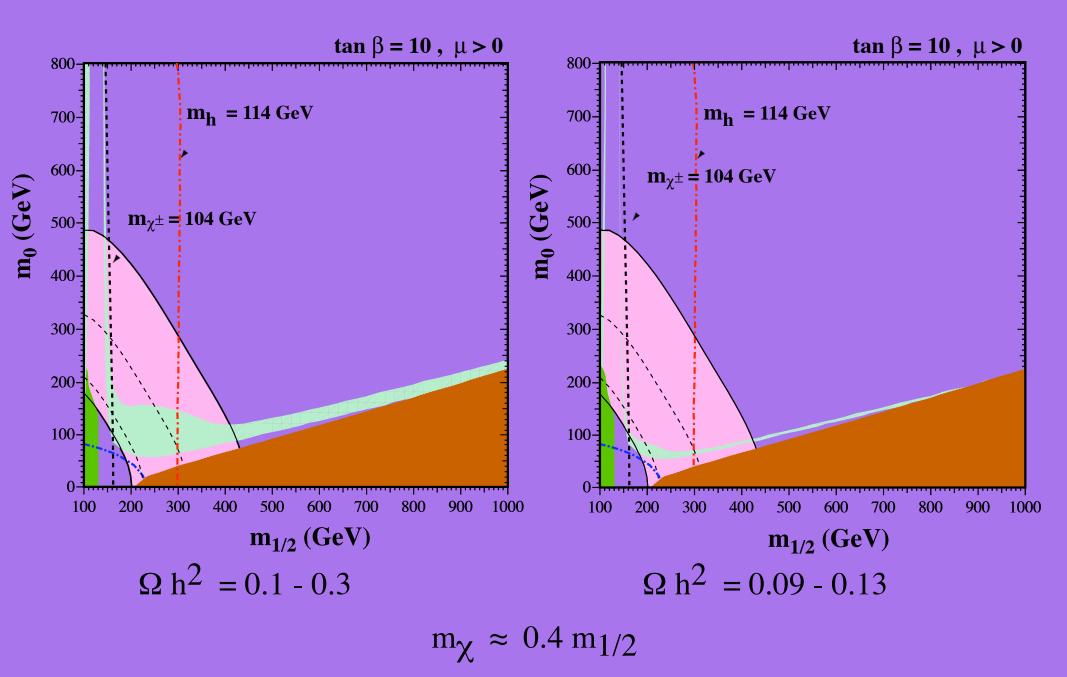
(in addition to m_0 , $m_{1/2}$, and A_0)

CMSSM conditions

• Instead CMSSM:

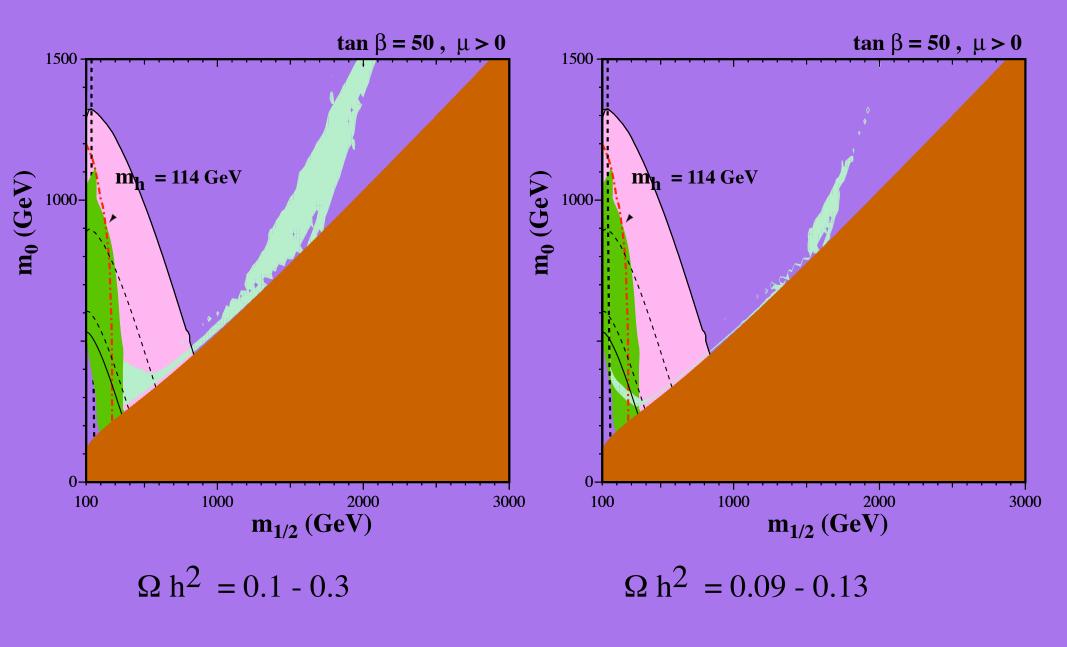
Input parameters: M_Z , m_1 , m_2 , $\tan \beta$ ($m_1 = m_2 = m_0$) predict μ , B, m_A

Effect of WMAP Densities



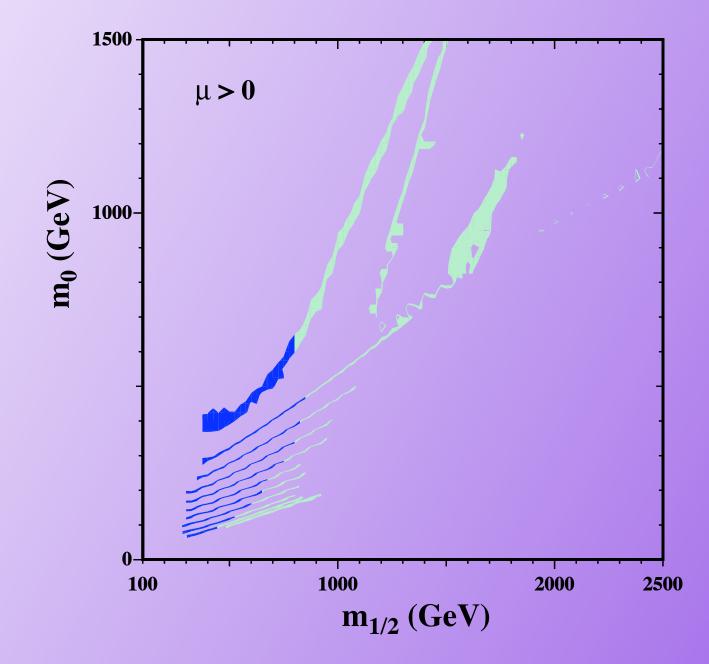
Ellis, Olive, Santoso, Spanos

Effect of WMAP Densities



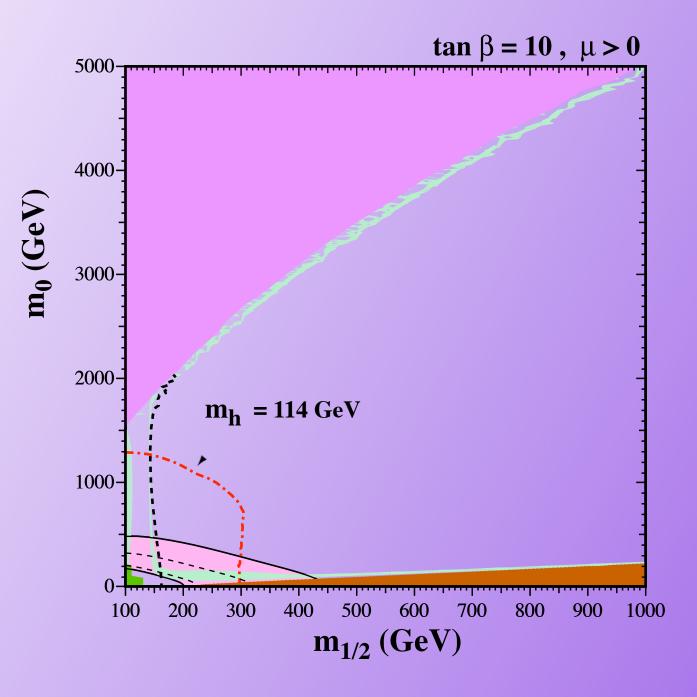
Ellis, Olive, Santoso, Spanos

Foliation in tan β



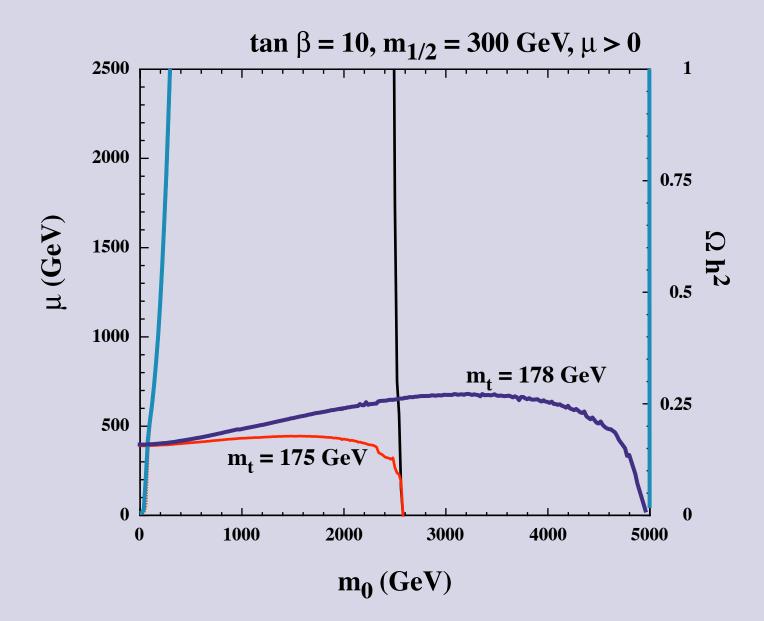
Focus Point Region

As m₀ gets very large, RGE's force μ to 0, allowing neutralino to become Higgsino like with an acceptable relic density.



Feng Matchev Moroi Wilczek

Large shift in focus point when $m_t = 178 \text{ GeV}$



Likelihood Analysis of the CMSSM parameter space Ellis, KAO, Santoso, Spanos

Includes

• Likelihood for the direct LEP Higgs Search

+ Global fit to precision electrowweak data

• Likelihood for b to s γ

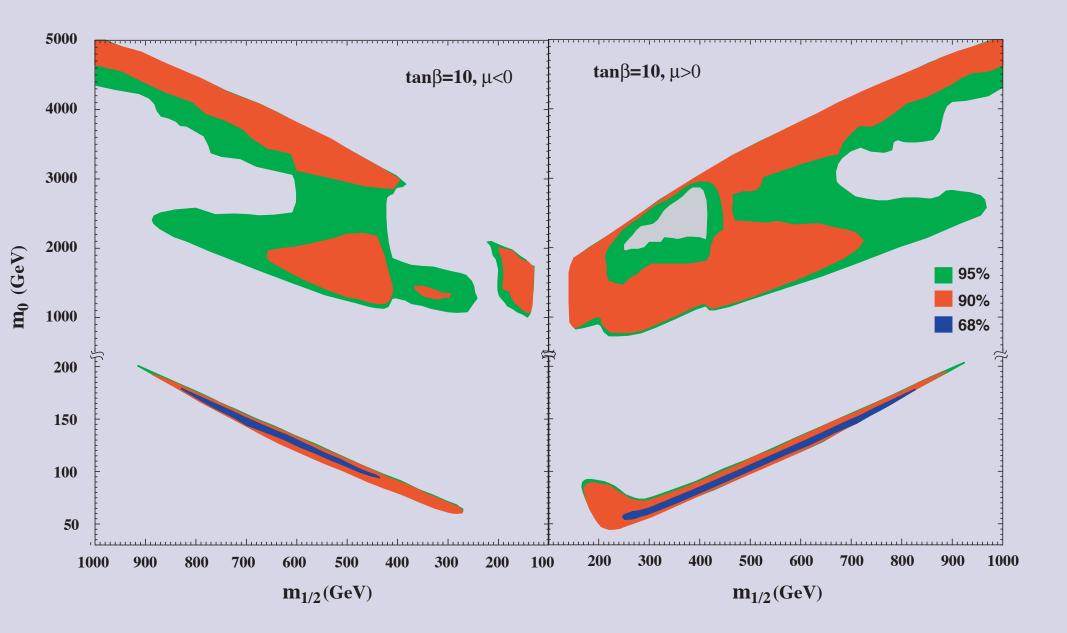
Gambino and Ganis

• g-2 data (optional)

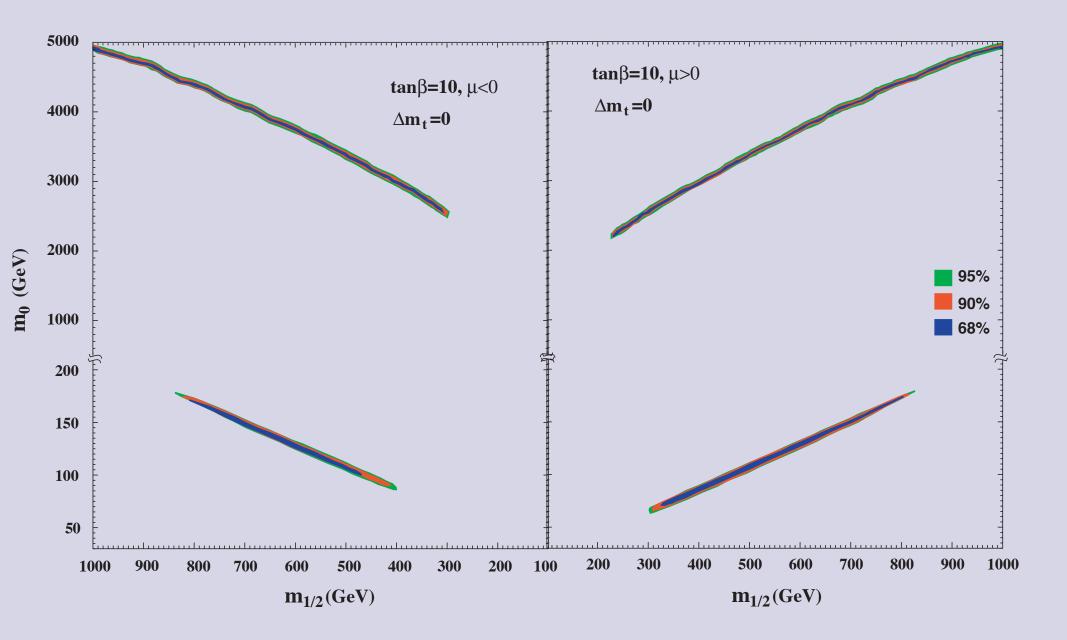
Davier et al

• Relic Density

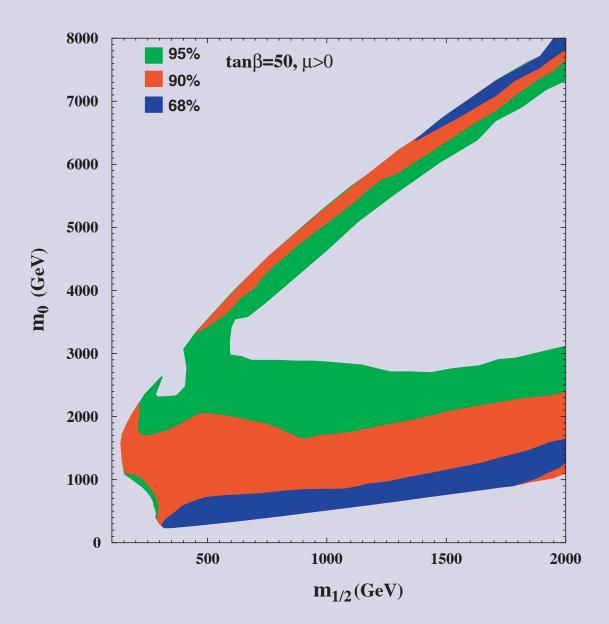
Likelihood Projections



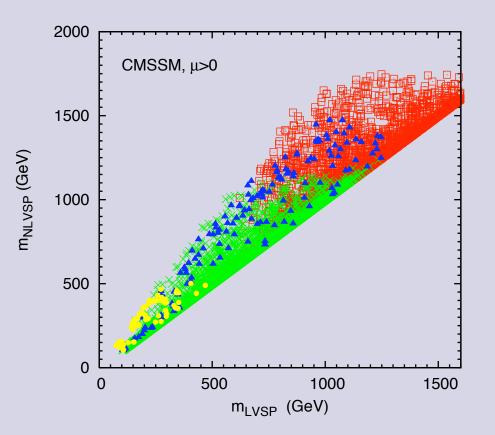
Likelihood Projections

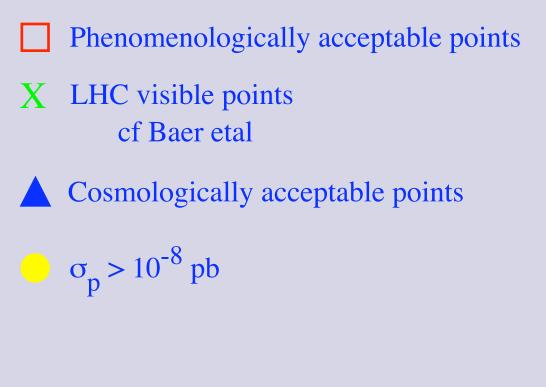


Likelihood Projections

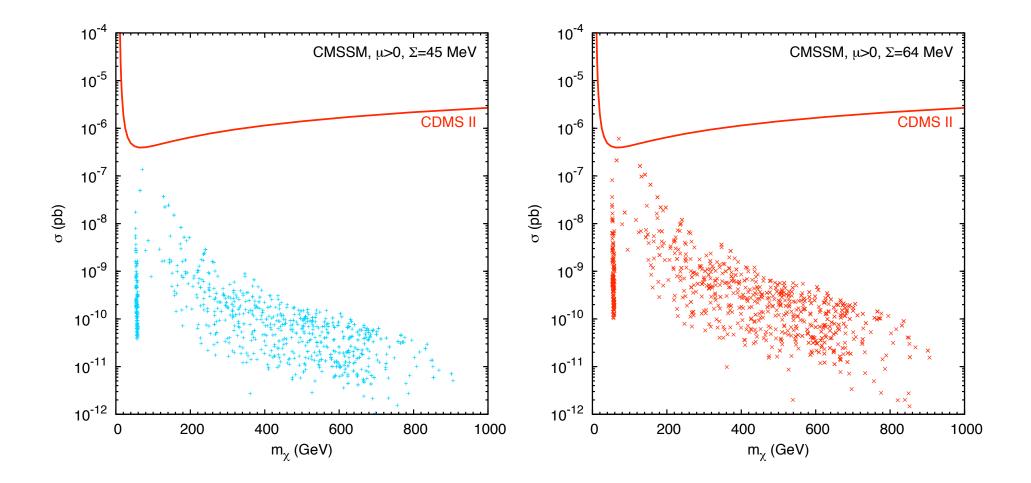


Visible Particle Masses

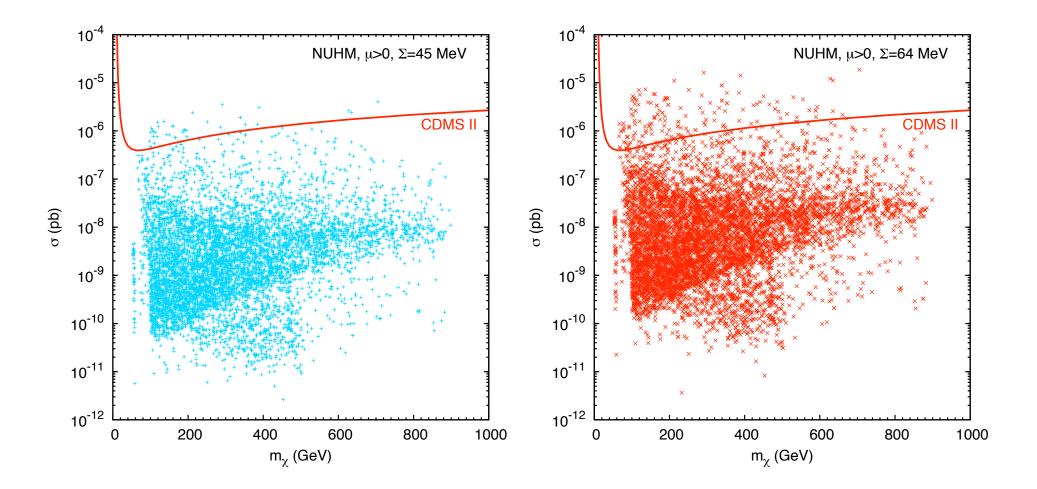




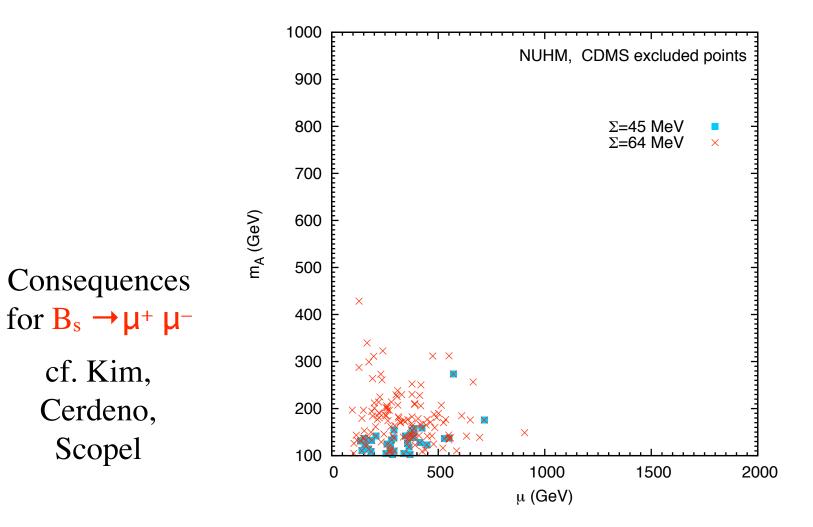
Direct Detection in the CMSSM



Direct Detection in the NUHM



CDMS Excluded models



Ellis, KAO, Santoso, Spanos

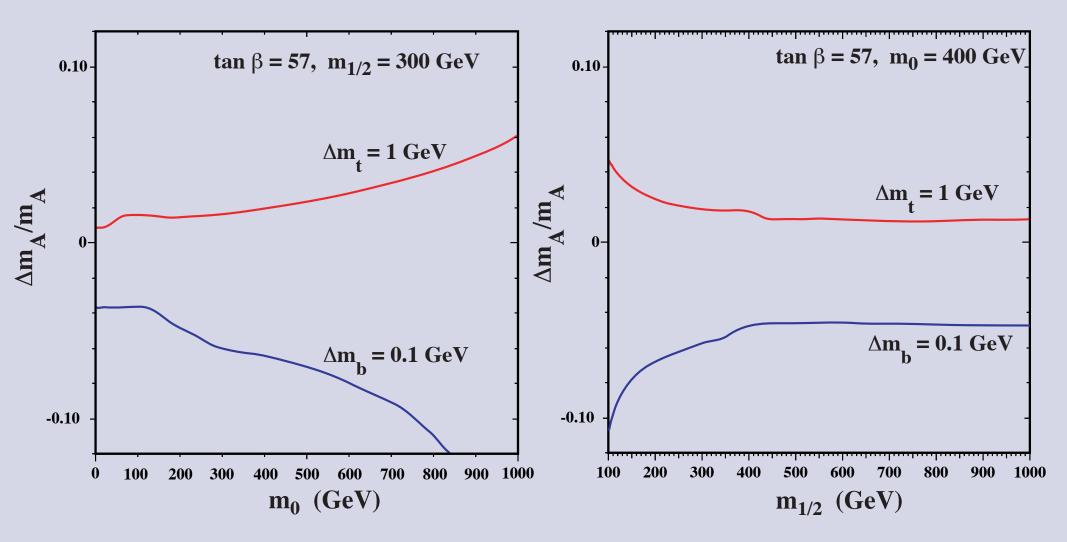
Uncertainties in $B_s \rightarrow \mu^+ \mu^-$

$$\mathcal{B}(B_s \to \mu^+ \,\mu^-) = \frac{G_F^2 \alpha^2}{16\pi^3} \frac{M_{B_s}^5 f_{B_s}^2 \tau_B}{4} |V_{tb} V_{ts}^*|^2 \sqrt{1 - \frac{4m_{\mu}^2}{M_{B_s}^2}} \\ \times \left\{ \left(1 - \frac{4m_{\mu}^2}{M_{B_s}^2} \right) |C_S|^2 + \left| C_P - 2 C_A \frac{m_{\mu}}{M_{B_s}^2} \right|^2 \right\} \,,$$

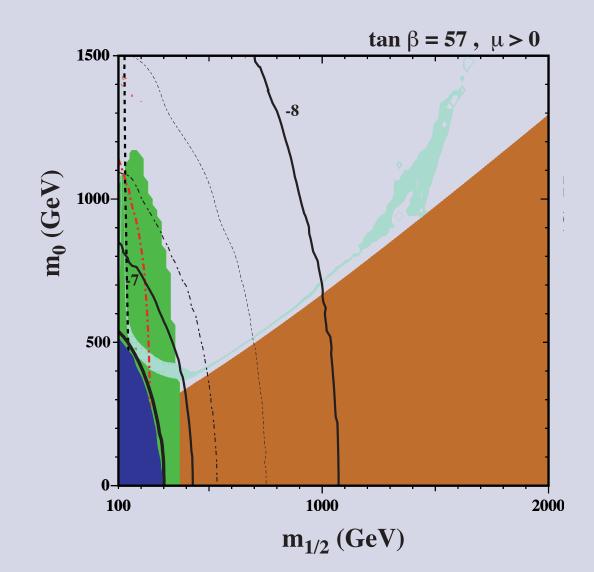
$$C_{S,P}^{\rm CT} = \mp \frac{m_{\mu} \tan^3 \beta}{\sqrt{2} M_W^2 m_A^2} \sum_{i=1}^2 \sum_{a=1}^6 \sum_{m,n=1}^3 [m_{\tilde{\chi}_i^{\pm}} D_3(y_{ai}) U_{i2}(\Gamma^{U_L})_{am} \Gamma_{imn}^a],$$

- B_s meson parameters f_B , m_B , τ_B
- $m_t = 178 \pm 4 \text{ GeV}$
- $m_b = 4.25 \pm 0.11 \text{ GeV}$

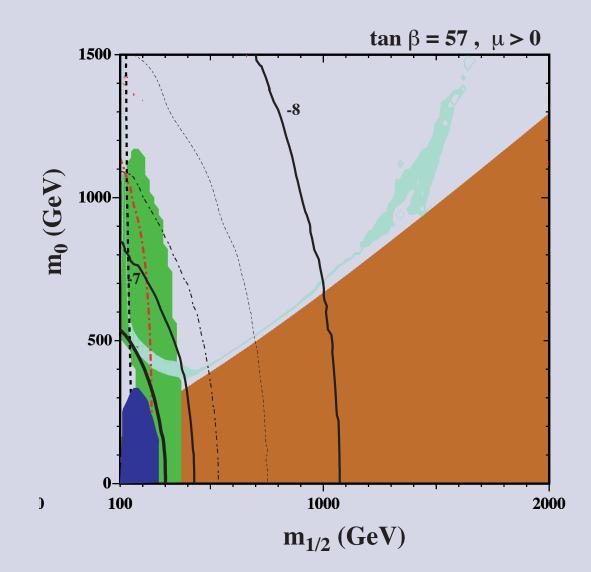
Sensitivity of m_A to m_t and m_b



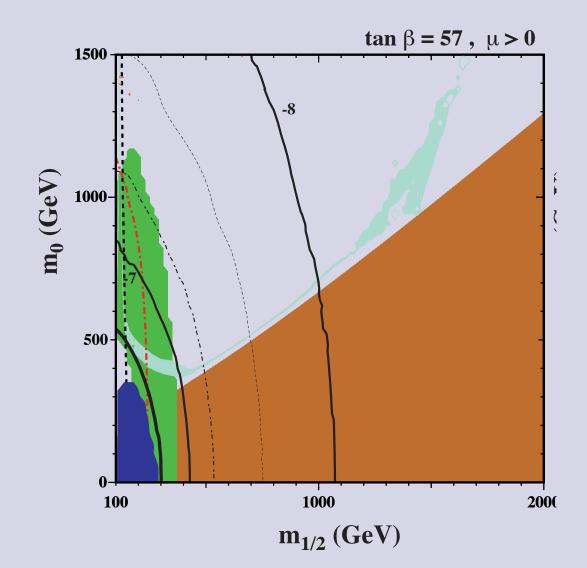
Effect of B parameters

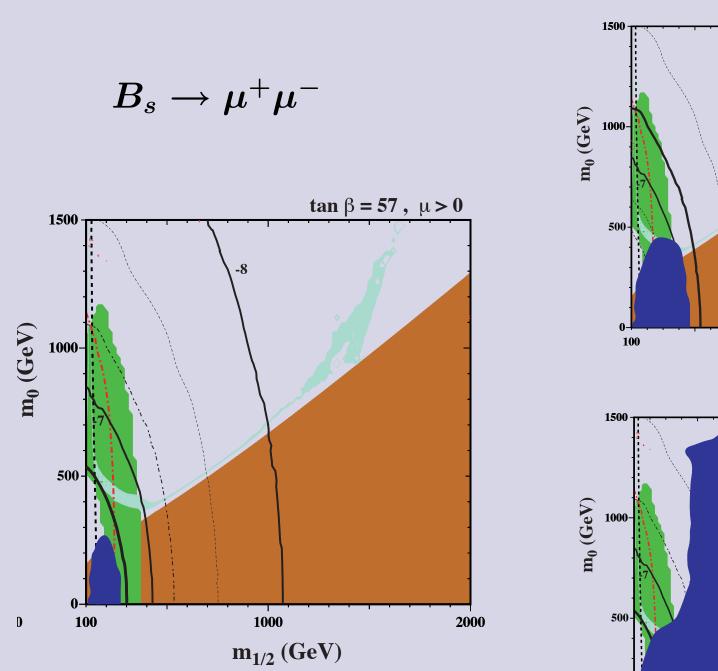


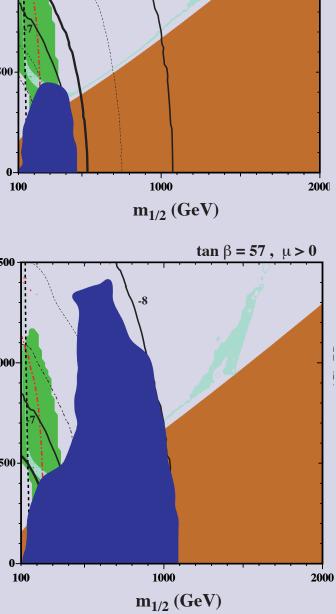
Effect of uncertainty in m_t



Effect of uncertainty in mb







-8

 $\tan\beta=57\ ,\ \mu>0$

Minimal Supergravity Models

e.g. Barbieri, Ferrara, Savoy

Nilles, Srednicki, Wyler

 $G = \phi \phi * + z z^* + \ln |W|^2; W = f(z) + g(\phi)$

$$V = \left| \frac{\partial g}{\partial \varphi} \right|^2 + m_{3/2} \left(\varphi \frac{\partial g}{\partial \varphi} + \left(\sqrt{3} \langle z \rangle - 3 \right) g + h.c. \right) + m_{3/2}^2 \varphi \varphi^*$$

3g for trilinear terms

3g for trilinear terms 2g for bilinear terms

For Polonyi models $\langle z \rangle = \sqrt{3} - 1$, and

 $m_0 = m_{3/2}$; $A_0 = (3 - \sqrt{3}) m_0$; $B_0 = (A_0 - 1) m_0$

Boundary conditions

• Input parameters: μ , m_1 , m_2 , B. predict M_7 , tan β , m_A

(in addition to m_0 , $m_{1/2}$, and A_0)

CMSSM conditions

Instead CMSSM:

Input parameters: M_7 , m_1 , m_2 , $\tan \beta$ ($m_1 = m_2 = m_0$) predict μ , B, m_{Δ} Very CMSSM conditions

• Then:

Input parameters: M_Z, m₁, m₂, B

predict μ , m_A, (an β)

Electroweak Symmetry Breaking conditions:

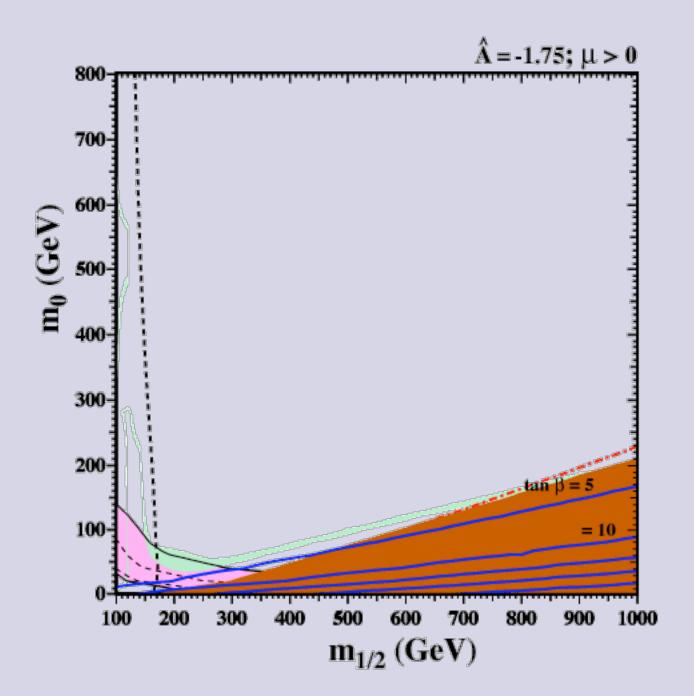
$$\mu^{2} = \frac{m_{1}^{2} - m_{2}^{2} \tan^{2} \beta + \frac{1}{2} m_{Z}^{2} (1 - \tan^{2} \beta) + \Delta_{\mu}^{(1)}}{\tan^{2} \beta - 1 + \Delta_{\mu}^{(2)}}$$

$$B\mu = -\frac{1}{2}(m_1^2 + m_2^2 + 2\mu^2)\sin 2\beta + \Delta_B$$

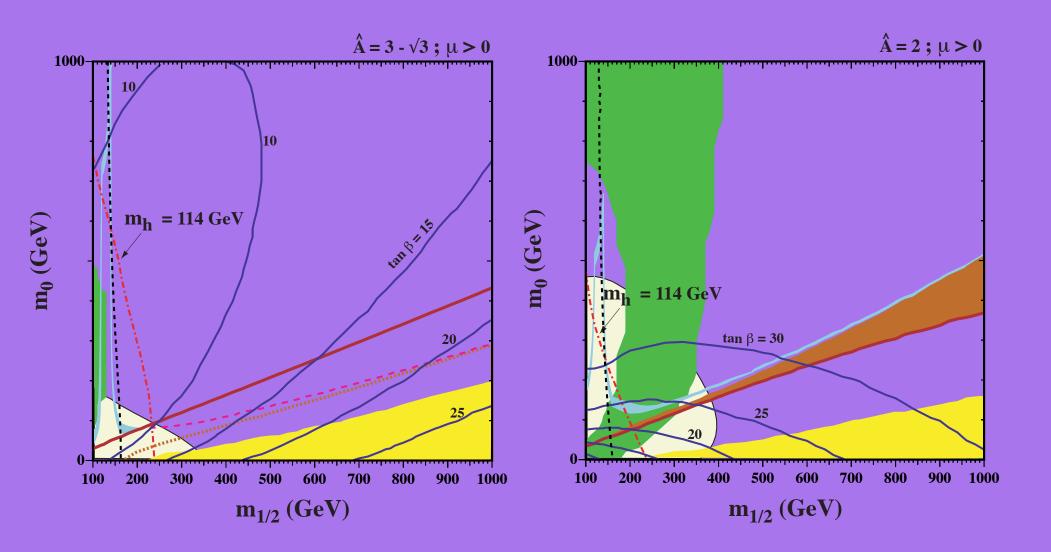
But Δ_{μ} depends on tan β and Δ_{B} depends on μ and tan β so one can not write down and expression for

$$\tan\beta = \dots$$

Of course it can be solved numerically



VCMSSM



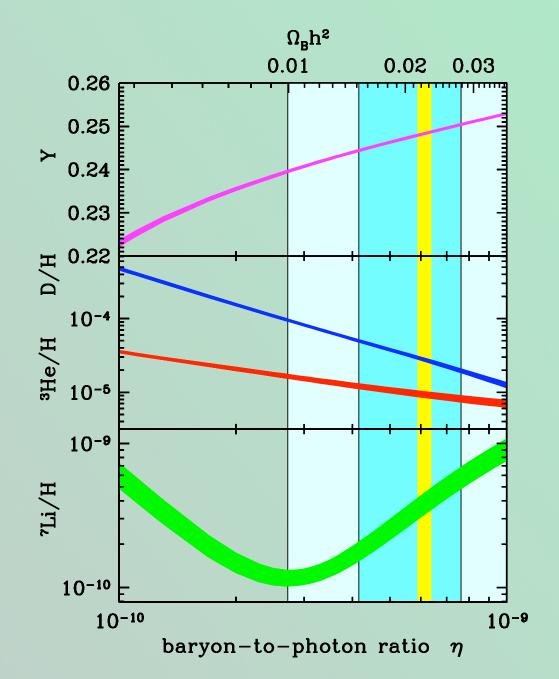
Gravitinos

- Strong cosmological limits on mass
 - $m_{3/2} < 1 \text{ keV} (\text{stable})$
 - or $m_{3/2} > 10$ TeV (unstable)
- Abundance limits after inflation (unstable)

$$egin{aligned} & au \sim 3 imes 10^8 \left(rac{100 GeV}{m_{3/2}}
ight)^3 \ & \quad rac{n_{3/2}}{n_\gamma} \equiv Y_{3/2} < 5 imes 10^{-14} \left(rac{100 GeV}{m_{3/2}}
ight) \qquad au \sim 10^8 s \end{aligned}$$

BBN Concordance

- Concordance rests on balance between interaction rates and expansion rate.
- Allows one to set constraints on:
 - Particle Types
 - Particle Interactions
 - Particle Masses
 - Fundamental Parameters



Cyburt, Fields, KAO

Limits on Unstable particles due to

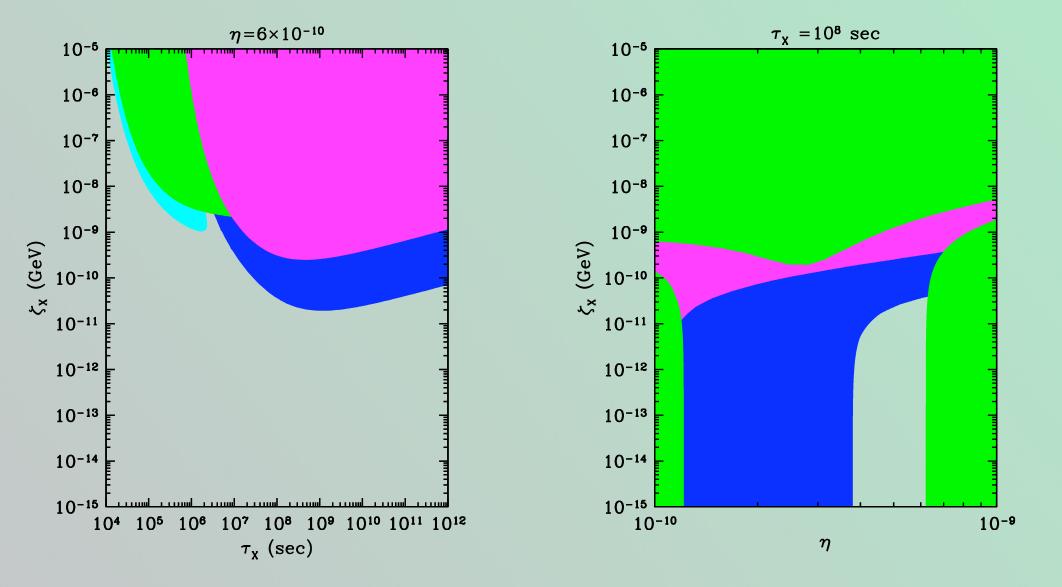
Photo-Destruction and -Production of Nuclei

3 free parameters

$$\zeta_X \equiv \frac{n_X^0}{n_\gamma^0} M_X = r M_X = 2r E_0,$$

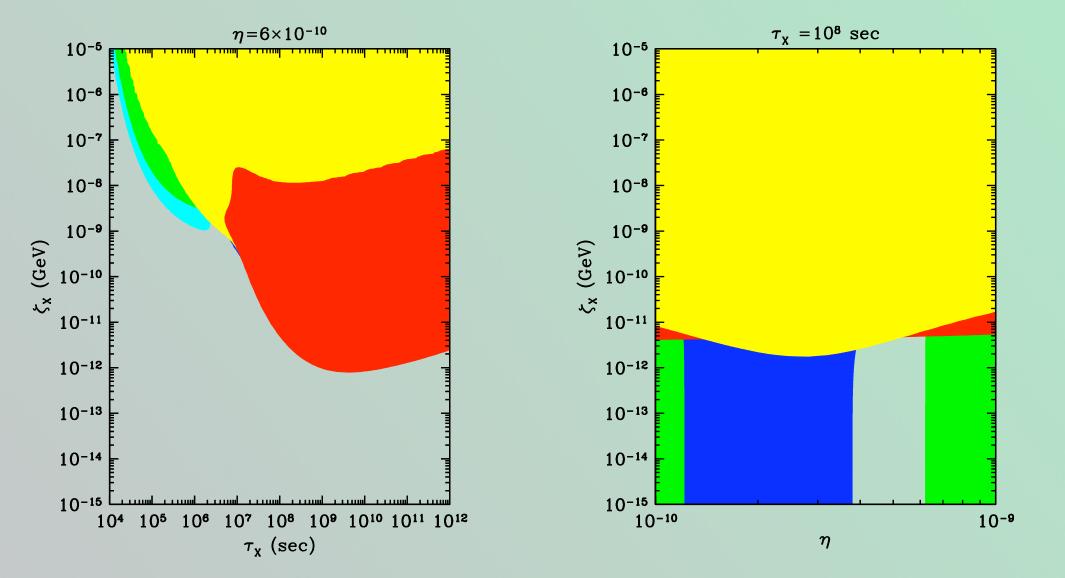
$$\tau_X \text{ and } \eta.$$

D, He and Li

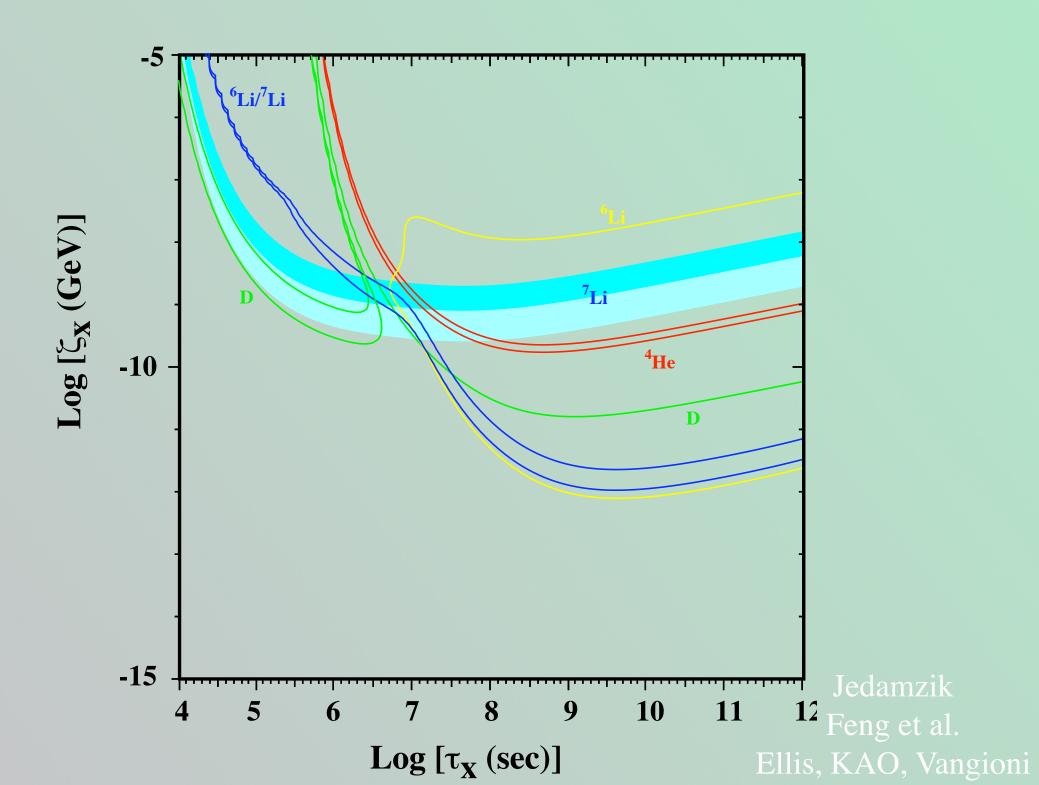


Cyburt, Ellis, Fields, KAO

"All"



Cyburt, Ellis, Fields, KAO



Gravitino Dark Matter and NSP decays

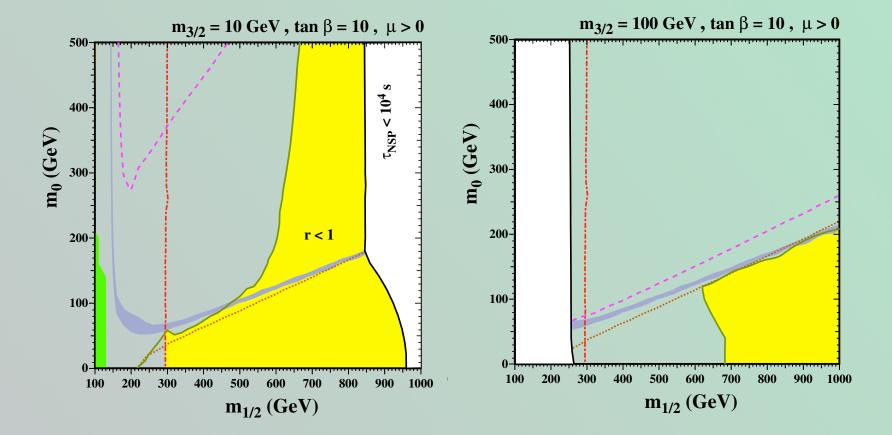
Ellis, Olive, Santoso, Spanos

 χ is now the NSP and can decay to gravitinos

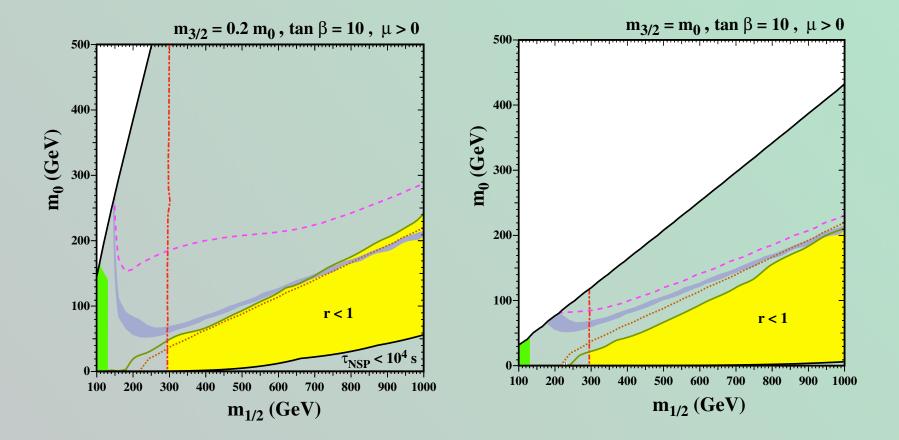
$$\Omega_{3/2} = \Omega_{\chi}(m_{3/2}/m_{\chi})$$
 $\Omega_X^0 h^2 = 3.9 \times 10^7 \text{ GeV}^{-1} \zeta_X$

There could be additional sources of gravitino dark matter (e.g. inflation)

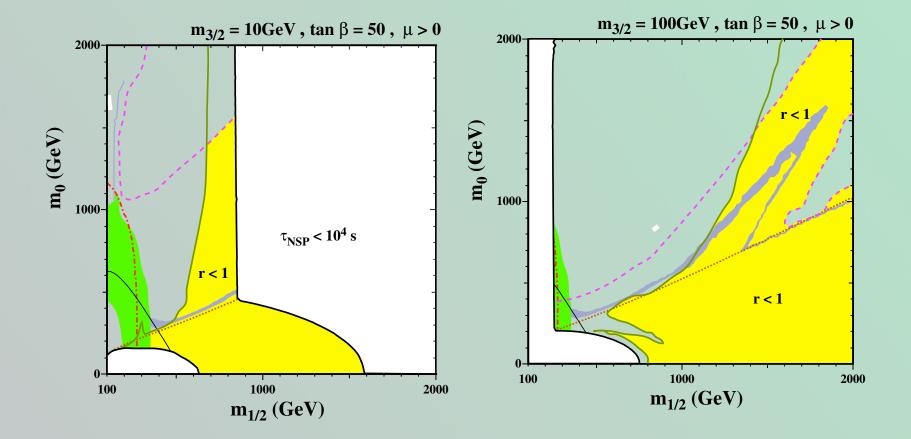
The CMSSM plane with gravitino dark matter



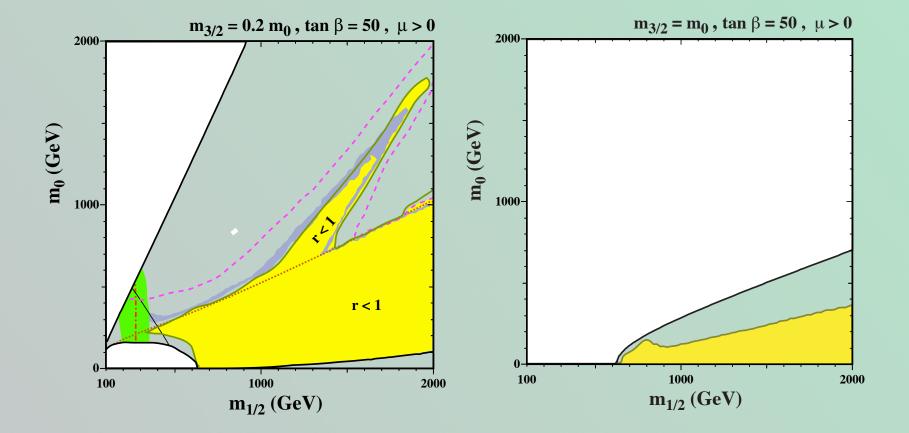
The CMSSM plane with gravitino dark matter



The CMSSM plane with gravitino dark matter

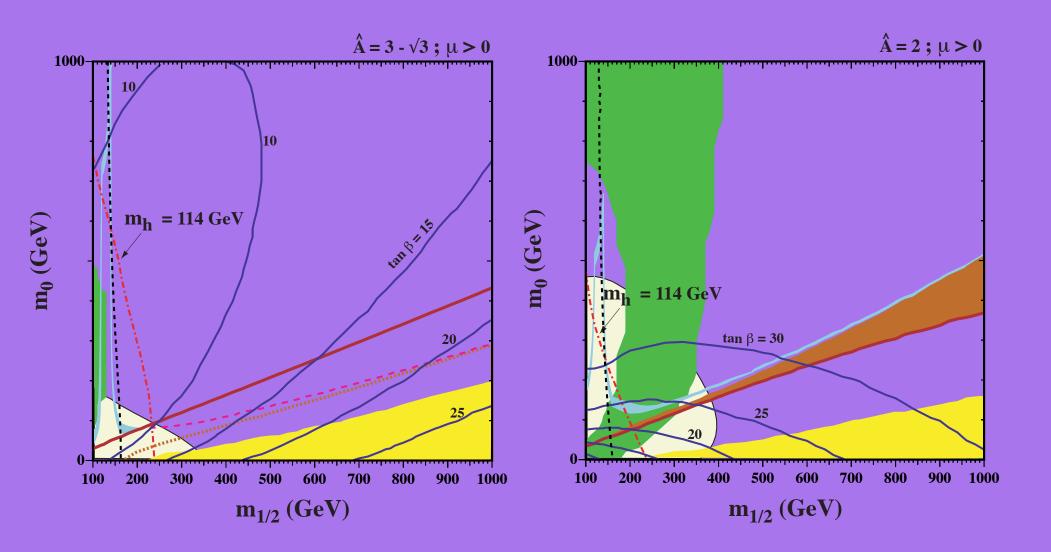


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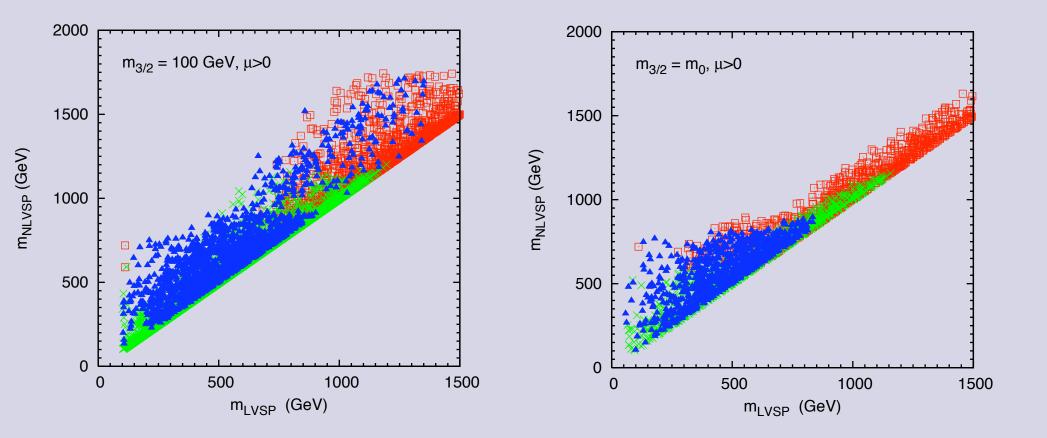


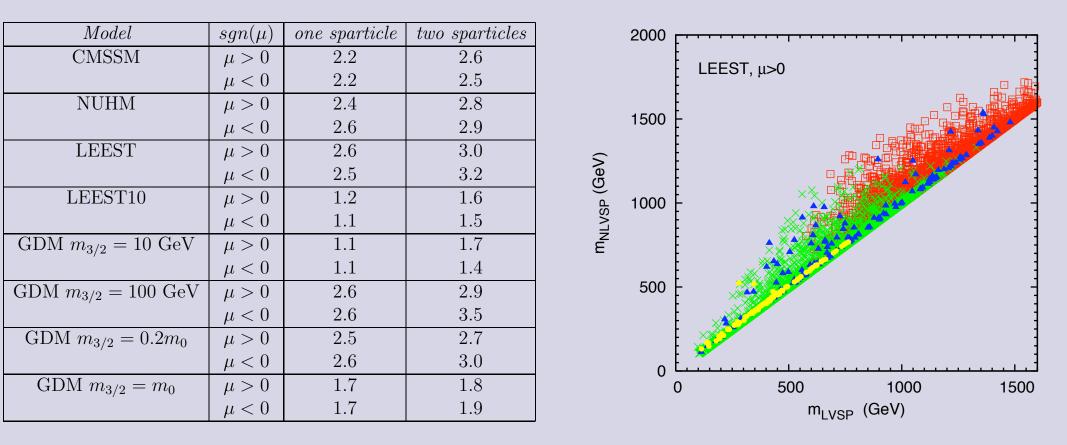
Ellis, KAO, Santoso, Spanos

VCMSSM



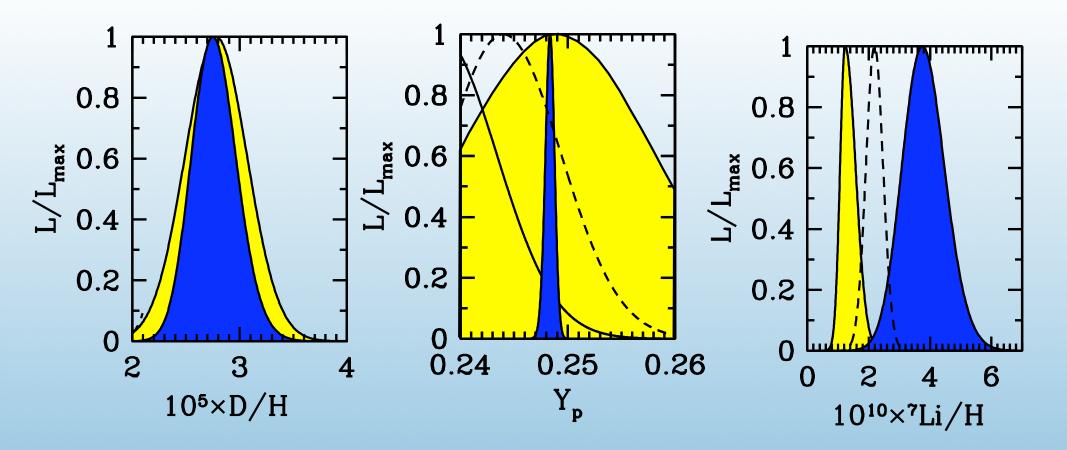
Visible Particle Masses





Required Center of Mass Energies

In most cases the LVSP is the τ but often there is a lower threshold due to the associated production of $\chi \chi_2$



Concordance of BBN with WMAP Cyburt, Fields, KAO

Possible sources for the Li discrepancy

- Stellar Depletion
 - lack of dispersion in the data, ⁶Li abundance
 - standard models (< .05 dex), models (0.2 0.4 dex)
- Nuclear Rates

Vauclaire & Charbonnel Pinsonneault et al.

- Restricted by solar neutrino flux

Coc et al. Cyburt, Fields, KAO

• Stellar parameters

 $rac{dLi}{dlng}=rac{.09}{.5} \qquad \qquad rac{dLi}{dT}=rac{.08}{100K}$

Possible sources for the Li discrepancy

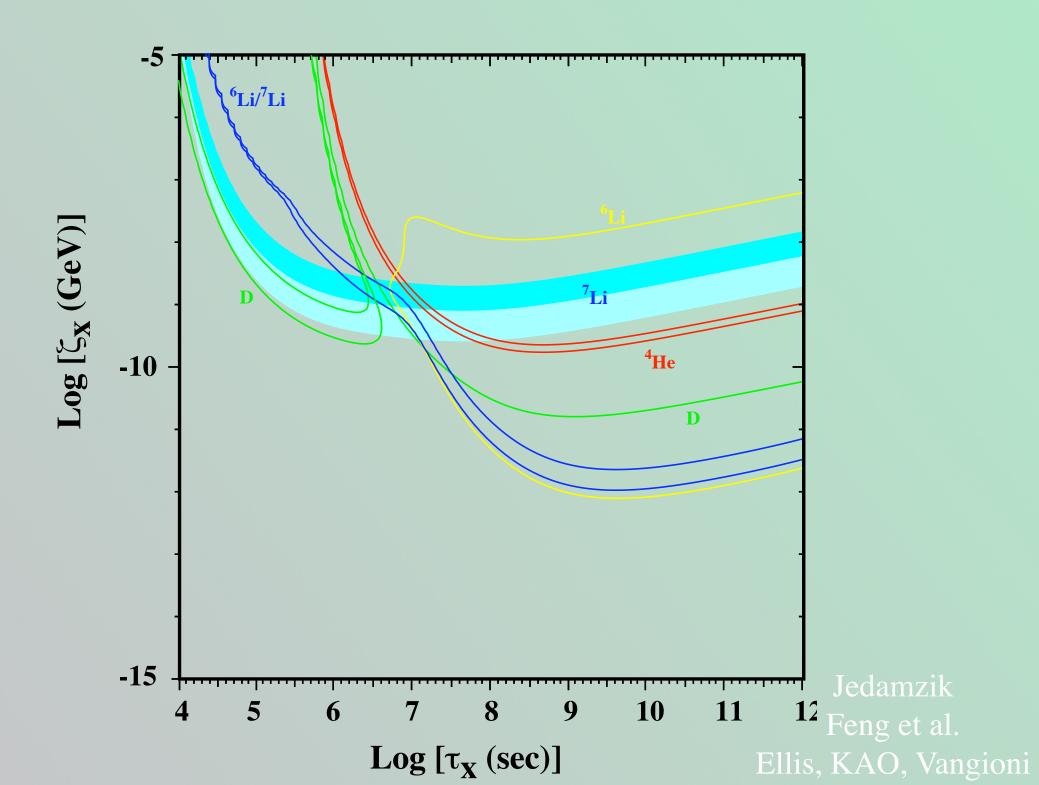
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Coc et al. Cyburt, Fields, KAO

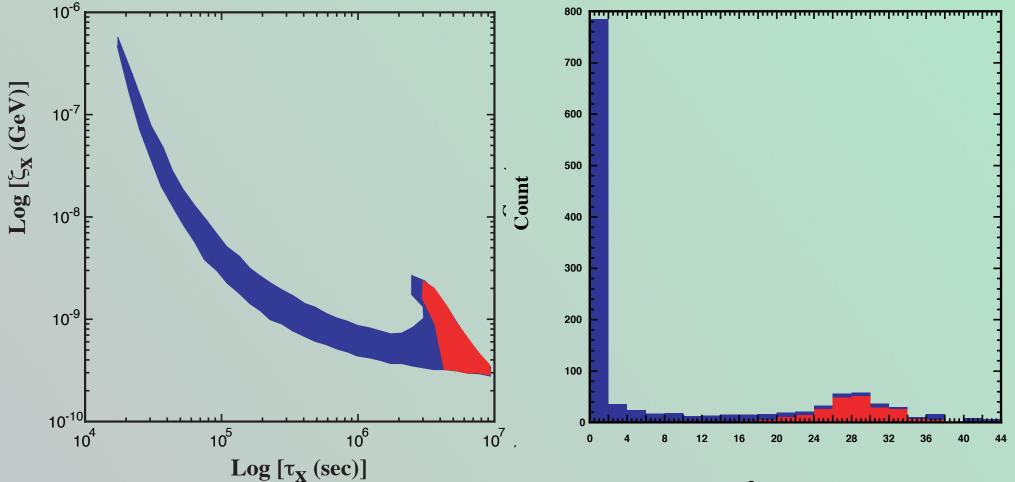
• Stellar parameters

 $\frac{dLi}{dlng} = \frac{.09}{.5} \qquad \qquad \frac{dLi}{dT} = \frac{.08}{100K}$

• Particle Decays



Blue: $D/H > 1.3 \times 10^{-5}$ Red: $D/H > 2.2 \times 10^{-5}$

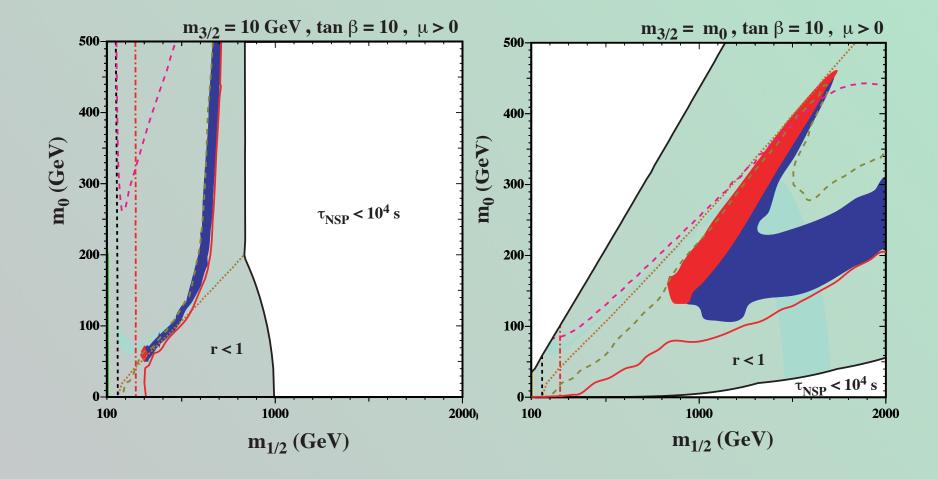


³He/D

Require $^{3}\text{He/D} < 1$

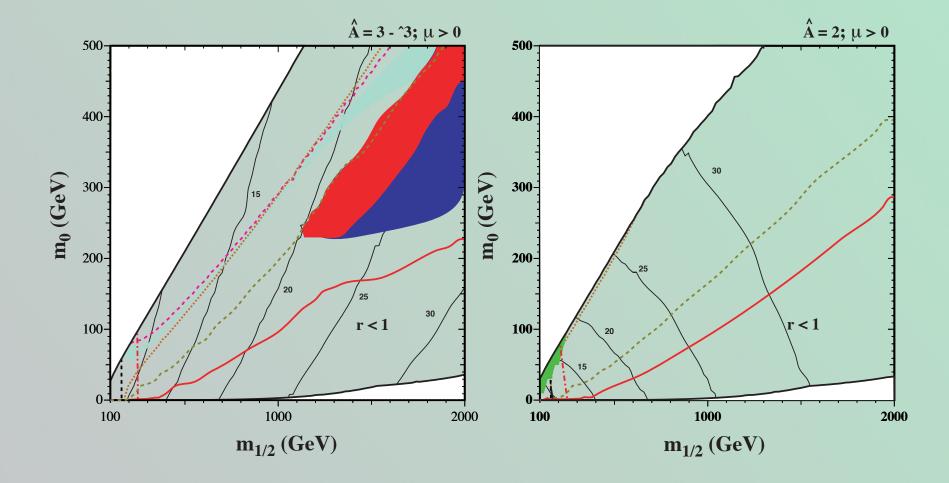
Ellis, KAO, Vangioni

Regions in the plane with Li/H < 3 x 10⁻¹⁰ Blue: D/H > 1.3 x 10⁻⁵ Red: D/H > 2.2 x 10⁻⁵

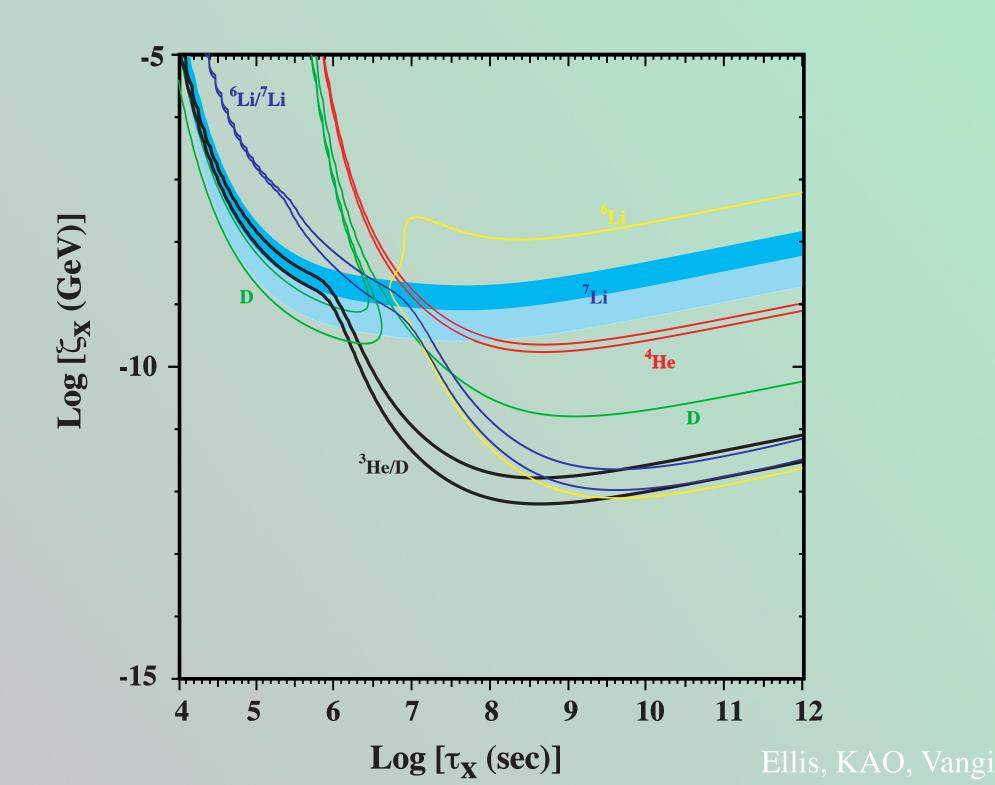


Ellis, KAO, Vangioni

Regions in the plane with Li/H < 3 x 10⁻¹⁰ Blue: D/H > 1.3 x 10⁻⁵ Red: D/H > 2.2 x 10⁻⁵



Ellis, KAO, Vangioni



Summary

• CMSSM Dark Matter:

Bulk regions; stau-coannihilation; A-pole funnel; focus point

- Spectrum Uncertainties: Limits excludable regions;
 cf: B_s → µ⁺ µ⁻
- VCMSSM ($B_0 = A_0 m_0$) ($m_0 = m_{3/2}$)
 - Restricted range for tan β for a given A₀.
 - No funnels or focus point
 - much of the parameter space predicts gravitino dark matter
- Gravitino Dark Matter
 - More parameter space allowed (including stau NLSP)
 - * Can not resolve Li problem